

Noise robustness and threshold of many-body quantum magic

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Outline

- Background and Preliminaries
- Noise robustness and threshold of magic
- Result 1. Magic induced by broader interactions is less robust
- Result 2. Local magic and magic threshold
- Outlook

Background

- Why quantum computing is potentially more powerful than classical computing?
 - "Magic", a source of quantum computational advantages
- Noise can significantly undermine the resource features of quantum systems.
 - Noisy intermediate-scale quantum (NISQ) technologies
 - Quantum error correction
 - Has been investigated for features like computational supremacy, entanglement ...

Aharonov D, Gao X, Landau Z, et al. A polynomial-time classical algorithm for noisy random circuit sampling[C]//STOC. 2023: 945-957.

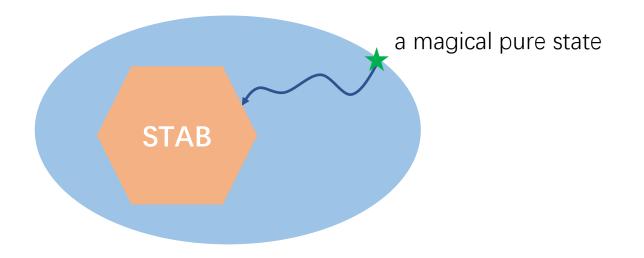
Background

• How noise effects impact magic?



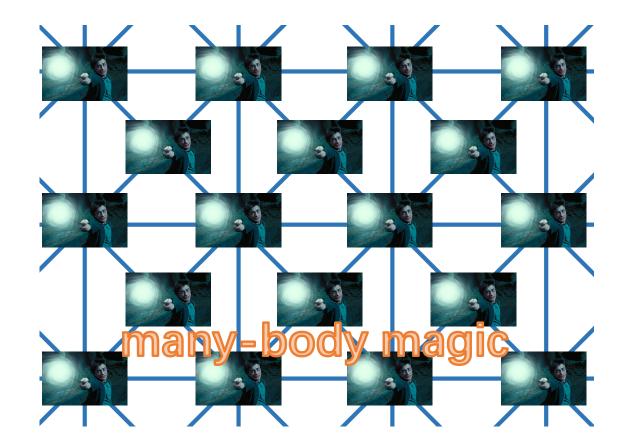
Background

- How noise effects impact magic?
- Noise drives pure states to mixed ones, and as it intensifies, the magic gradually decays and eventually vanishes at a certain point as the state is brought inside the stabilizer hull.



- How noise effects impact magic in large systems with different entanglement structures?
- Generates insights into
 - Interplay between magic and entanglement
 - Design of circuits
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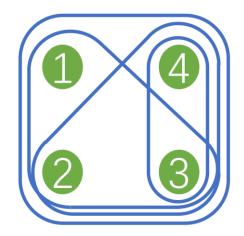
- A neat yet highly versatile model for entangled magic states: hypergraph states.
 - many-body physics
 - measurement-based quantum computing (MBQC)
- Also provide an apt playground for concretely investigating the relation between entanglement structures and magic properties

Levin M, Gu Z C. Braiding statistics approach to symmetry-protected topological phases[J]. PRB, 2012, 86(11): 115109. Miller J, Miyake A. Latent computational complexity of symmetry-protected topological order with fractional symmetry[J]. PRL, 2018, 120(17): 170503.

Raussendorf R, Briegel H J. A one-way quantum computer[J]. PRL, 2001, 86(22): 5188. Miller J, Miyake A. Hierarchy of universal entanglement in 2D measurement-based quantum computation[J]. npj Quantum Information, 2016, 2(1): 1-6.

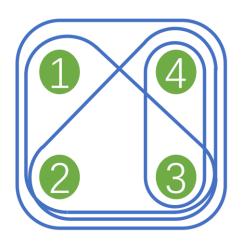
Liu Z W, Winter A. Many-body quantum magic[J]. PRX Quantum, 2022, 3(2): 020333.

• A neat yet highly versatile model for entangled magic states: hypergraph states.



$$E = \{\{3, 4\}, \{1, 2, 3\}, \{2, 3, 4\}, \{1, 2, 3, 4\}\}$$
$$|\Psi_G\rangle = CZ_{34}CCZ_{123}CCZ_{234}CCCZ_{1234}|+^4\rangle$$

• A neat yet highly versatile model for entangled magic states: hypergraph states.



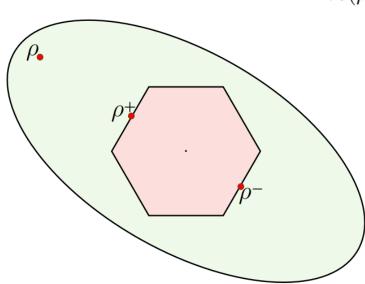
Let $C^{n-1}Z = \text{diag}(1, \dots, 1, -1)$ denote the multicontrolled-Z gate on n-qubits, with $C^0Z = Z$. Given a hypergraph $G = \{[n], E\}$, where $[n] := \{1, \dots, n\}$ is the set of vertices and $E \subset 2^{[n]}$ is the set of hyperedges, an associated hypergraph state is defined by

$$|\Psi_G\rangle = \prod_{e \in E} \mathcal{C}^{|e|-1} Z_e |+^n\rangle,$$

where $C^{|e|-1}Z_e$ is the multi-controlled-Z gate applied to the qubits in $e \subset [n]$.

The number of vertices contained in an edge e is referred to as the *degree* of the edge.

• Magic measure for mixed states: robustness of magic (RoM)



$$\mathcal{R}(\rho) = \min_{\rho^+, \rho^- \in \text{STAB}_n} \left\{ 2a + 1 \mid \rho = (a+1)\rho^+ - a\rho^-, a \ge 0 \right\}$$

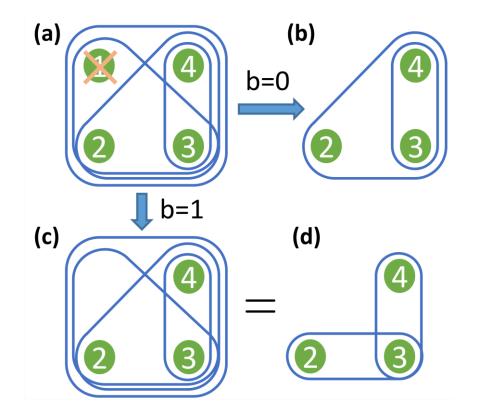
- 1. $\mathcal{R}(\cdot)$ is faithful, i.e., $\mathcal{R}(\sigma) = 1$ iff $\sigma \in \text{STAB}$;
- 2. For all trace-preserving stabilizer channels \mathcal{E} , we have $\mathcal{R}(\mathcal{E}(\rho)) \leq \mathcal{R}(\rho)$;
- 3. $\mathcal{R}(\sigma \otimes \rho) = \mathcal{R}(\rho)$ for $\sigma \in \text{STAB}$;
- 4. For a set of states $\{\rho_k\}$ and a set of real numbers $\{p_k\}$ satisfying $\sum_k p_k = 1$, the *convexity* of $\mathcal{R}(\cdot)$ implies that $\mathcal{R}(\sum_k p_k \rho_k) \leq \sum_k |p_k| \mathcal{R}(\rho_k);$

5. Classical simulation overhead: $\mathcal{R}(\rho)^2$.

Howard M, Campbell E. Application of a resource theory for magic states to fault-tolerant quantum computing[J]. PRL, 2017, 118(9): 090501.

• Consider the n-qubit independent depolarizing noise

 $\mathcal{E}_{\lambda}^{\otimes n} = ((1-\lambda)\mathcal{I} + \lambda\mathcal{G})^{\otimes n}, \text{ where } \mathcal{G}(\sigma) = \operatorname{Tr}(\sigma)\mathbb{I}_2/2.$



For an *n*-qubit hypergraph state $\Psi := |\Psi\rangle\langle\Psi|$, and a subset of qubits $I \subset [n]$, consider tracing out the qubits in *I*. We have

$$\operatorname{Tr}_{I}(\Psi) = \frac{1}{2^{|I|}} \sum_{\mathbf{b} \in \mathbb{Z}_{2}^{|I|}} \Psi^{(I,\mathbf{b})}, \qquad (1)$$

where $\Psi^{(I,\mathbf{b})}$ are (n - |I|)-qubit hypergraph states obtained by removing vertices and edges from Ψ .

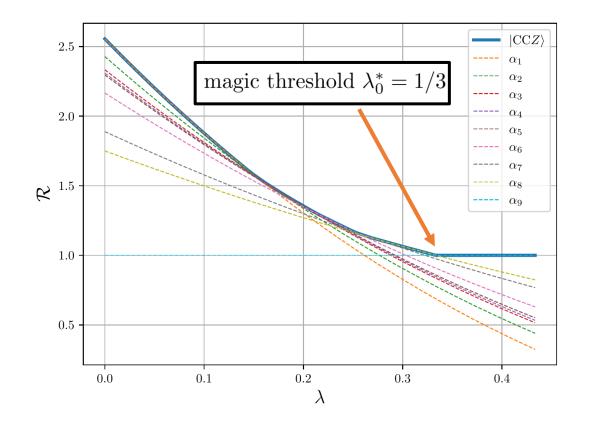
- The decay profile of magic: $\mathcal{R}\left(\mathcal{E}_{\lambda}^{\otimes n}(
 ho)
 ight)$
- The magic noise threshold above which the magic is eliminated.

Definition 1. For a state ρ and $\epsilon \geq 0$, we call $\lambda_{\epsilon}^*(\rho) := \inf_{\mathcal{R}(\mathcal{E}_{\lambda}^{\otimes n}(\rho)) \leq 1+\epsilon} \lambda$ the ϵ -magic noise threshold. For a family of states $\{\rho_n\}$ where ρ_n is an *n*-qubit state, if for a fixed $\epsilon \geq 0$ we have $\liminf_{n \to \infty} \lambda_{\epsilon}^*(\rho_n) > 0$, we say $\{\rho_n\}$ has a non-vanishing magic threshold.

• Warm up example: $\Phi = |CCZ\rangle\langle CCZ|$, where $|CCZ\rangle \coloneqq CCZ|+^3\rangle$

 $\mathcal{R}(\mathcal{E}_{\lambda}^{\otimes n}(\Phi)) = \max_{1 \leq j \leq 9} \alpha_j(\lambda), \text{ where } \alpha_j(\lambda) = \operatorname{Tr}\left(\mathcal{E}_{\lambda}^{\otimes 3}(\Phi)A_j\right) \text{ are polynomials in } \lambda \text{ with degree at most } 3, \text{ with } A_j \text{ some Hermitian matrices.}$

• Computing the decay profile for n > 5 is hard in general.



- Key technique: upper and lower bounds for noisy RoM.
- Upper bound: by convexity of $\mathcal{R}(\cdot)$

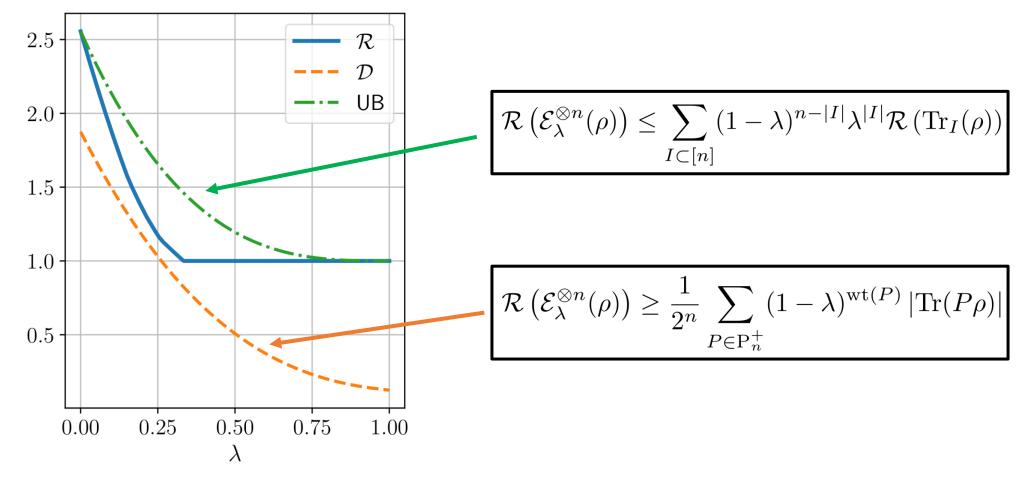
$$\mathcal{R}\left(\mathcal{E}_{\lambda}^{\otimes n}(\rho)\right) \leq \sum_{I \subset [n]} (1-\lambda)^{n-|I|} \lambda^{|I|} \mathcal{R}\left(\mathrm{Tr}_{I}(\rho)\right)$$

• Lower bound: by $\mathcal{R}(\rho) \ge \mathcal{D}(\rho) := \frac{1}{2^n} \sum_{P \in \mathbf{P}_n^+} |\operatorname{Tr}(P\rho)|$

$$\mathcal{R}\left(\mathcal{E}_{\lambda}^{\otimes n}(\rho)\right) \geq \frac{1}{2^{n}} \sum_{P \in \mathcal{P}_{n}^{+}} (1-\lambda)^{\operatorname{wt}(P)} \left|\operatorname{Tr}(P\rho)\right|$$

Howard M, Campbell E. Application of a resource theory for magic states to fault-tolerant quantum computing[J]. PRLs, 2017, 118(9): 090501. Leone L, Oliviero S F E, Hamma A. Stabilizer rényi entropy[J]. PRL, 2022, 128(5): 050402.

$\rho = |CCZ\rangle\langle CCZ|$



Result 1. High-degree edges are fragile

- Hypergraph state $\Phi_n = |C^{n-1}Z\rangle \langle C^{n-1}Z|$, a single edge containing all qubits.
 - we can show that $\mathcal{R}(|C^{n-1}Z\rangle) < 5$, for all n
 - magic spreads more extensively, more exposed to noise as n increases

$$2\left(1-3\lambda/4\right)^n \lesssim \mathcal{R}\left(\mathcal{E}_{\lambda}^{\otimes n}\left(\Phi_n\right)\right) \le 1+4\left(1-\lambda/2\right)^n$$

Theorem 1. For a fixed noise rate $0 < \lambda < 1$ and any $\epsilon > 0$, the family $\{|C^{n-1}Z\rangle\}$ has magic threshold $\lambda_{\epsilon}^* = \Theta(n^{-1})$.

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• $C^{n-1}Z$ variants of magic state distillation and injection?

At least $\Omega((1 - \lambda/2)^{-n})$ copies of $\mathcal{E}_{\lambda}^{\otimes n}(\Phi_n)$ are required to achieve constant RoM, and therefore, to distill Φ_n with constant success probability.

Result 1. High-degree edges are fragile

• Furthermore, in an arbitrary hypergraph state Ψ , the magic contributed by high-degree edges is fragile to noise.

Theorem 2. Let Ψ be an n-qubit hypergraph state. Consider adding K edges $e_1, \dots, e_K \subset [n]$ to Ψ to obtain a new state Φ . Then $\mathcal{R}\left(\mathcal{E}_{\lambda}^{\otimes n}(\Phi)\right)$ is upper bounded by $\mathcal{R}\left(\mathcal{E}_{\lambda}^{\otimes n}\left(\Psi\right)\right) + C_{\Psi}\sum_{\emptyset \neq J \subset [K]} (5^{|J|} + 1)\left(1 - \frac{\lambda}{2}\right)^{|\cup_{j \in J} e_j|},$ where $C_{\Psi} := \max_{I \subset [n], \mathbf{s} \in \mathbb{Z}_2^{|I|}} \mathcal{R}(\Psi^{(I,\mathbf{s})})$ is a constant depend-ing on Ψ .

For an *n*-qubit state ρ , define its maximum K-local RoM by

$$\mathcal{M}_{K}(\rho) := \max_{J \subset [n], |J| = K} \mathcal{R}\left(\operatorname{Tr}_{\overline{J}}(\rho)\right),$$

where $\overline{J} := [n] - J$.

For a family of states $\{\rho_n\}$, if there exists a constant Ksuch that $\liminf_{n\to\infty} \mathcal{M}_K(\rho_n) > 1$ (or $\lim_{n\to\infty} \mathcal{M}_K(\rho_n) = 1$), we say $\{\rho_n\}$ has *non-vanishing* (or vanishing) local magic.

Proposition 3. Suppose $\{\rho_n\}$ is a family of states with non-vanishing local magic, then it has a non-vanishing magic threshold.

Let J_n $(|J_n| = K)$ denote the K qubits on which ρ_n has maximum RoM. Since partial trace commutes with local noises, to ensure $\mathcal{E}_{\lambda}^{\otimes n}(\rho_n) \in \text{STAB}_n$, we must have $\mathcal{E}_{\lambda}^{\otimes K}(\text{Tr}_{\overline{J_n}}(\rho_n)) \in \text{STAB}_K$. Since the K-qubit state $\text{Tr}_{\overline{J_n}}(\rho_n)$ has RoM \geq constant for all sufficiently large n, it can withstand constant rate of noises before falling into STAB_K .

Theorem 4. There exist efficiently preparable families of many-body states with vanishing local magic yet a non-vanishing magic threshold.

For example, consider the family of 3-complete hypergraph states $\{\Gamma_n\}$, where Γ_n corresponds to the hypergraph with *n* vertices and all possible edges of degree 3. Evidently, $\{\Gamma_n\}$ has gate complexity $\mathcal{O}(n^3)$.

Theorem 4. There exist efficiently preparable families of many-body states with vanishing local magic yet a non-vanishing magic threshold.

For any constant K, the reduced density matrix of Γ_n on K qubits is a convex combination of four (n-K)-qubit hypergraph states, with all four coefficients approaching 1/4 as $n \to \infty$, which can be proved to form a stabilizer state. Specifically, in Proposition 13, we show that $\mathcal{M}_K(\Gamma_n) \leq 1 + 2^{1+\frac{3}{2}K - \frac{n}{2}}$, which tends to 1 for all K as $n \to \infty$, indicating that $\{\Gamma_n\}$ has vanishing local magic.

Theorem 4. There exist efficiently preparable families of many-body states with vanishing local magic yet a non-vanishing magic threshold.

By upper and lower bounds for $\mathcal{R}\left(\mathcal{E}_{\lambda}^{\otimes n}\left(\Gamma_{n}\right)\right)$, the magic threshold for $\{\Gamma_{n}\}$ satisfies $0.39 \leq \lambda_{\epsilon}^{*} \leq 0.78$, demonstrating a non-vanishing threshold.

• Insights into interplay between magic and entanglement:

 Γ_n has a vanishing amount of magic even when restricted to $\Theta(n)$ qubits: dividing Γ_n evenly among 4 parties results in a vanishing amount of magic for each party, as

$$\mathcal{M}_{n/4}(\Gamma_n) \le 1 + 2^{1-n/8} \to 1.$$

However, the total magic remains substantial, as $\mathcal{R}(\Gamma_n) \geq \mathcal{D}(\Gamma_n) = \Theta(2^{n/2}).$

• Insights into interplay between magic and entanglement:

The "non-local magic" in these states is embedded in entanglement and can withstand a constant rate of noise. We anticipate that such states may find interesting applications through "magic hiding" or "magic secret sharing".

| Hypergraph state Ψ | Edge degree | Threshold | $\mathcal{R}(\mathcal{E}_{\lambda}^{\otimes n}(\Psi))$ upper bound | $\mathcal{R}(\mathcal{E}_{\lambda}^{\otimes n}(\Psi))$ lower bound | Local magic |
|-----------------------------------|-------------|---|--|--|--|
| $ \mathrm{CC}Z angle$ | | $\lambda_0^* = 1/3$ | $\mathcal{R}(\mathcal{E}_{\lambda}^{\otimes n}(\Psi)) = \max_{1 \leq i}$ | $_{j\leq 9} \alpha_j(\lambda)$ [Fig. 3] | 23/9 |
| Union Jack lattice | 3 | $\lambda_0^* \ge \text{constant} > 0$ | | — | 1.0078 |
| 3-complete hypergraph | | $0.39 \le \lambda_{\epsilon}^* \le 0.78$ | $1 + 2^{n+1} (1 - (1 - 2^{-3/2})\lambda)^n$ | $\approx 2^{(n-3)/2} (1 - 3\lambda/4)^n$ | $1 + \mathcal{O}(2^{-n/2})$ |
| 4-complete hypergraph | 4 | $\lambda_0^* \ge \text{constant} > 0$ | $\mathcal{O}ig((2-\lambda)^nig)$ | $\mathcal{D}\left(\mathcal{E}_{\lambda}^{\otimes n}(\Psi) ight)$ [Fig. 5] | 1.25 + o(1) |
| High-degree hypergraphs | $\geq m$ | $\lambda_{\epsilon}^* = \mathcal{O}(1/m)$ | $1 + \operatorname{poly}(m)(1 - \lambda/2)^m$ | _ | 1 + o(1) |
| $ \mathrm{C}^{n-1}Z angle$ | n | $\lambda_{\epsilon}^* = \Theta(1/n)$ | $1 + 4(1 - \lambda/2)^n$ | $\approx 2\left(1-3\lambda/4\right)^n$ | $1 + \Theta\left((1/2)^n\right)$ |
| Qudit $ \mathbf{C}^{n-1}Z\rangle$ | n | $\lambda_0^* \ge 0.42 \ (d=3)$ | $1 + 4M_d(d-1)^n \left(1 - \frac{d-1}{d}\lambda\right)^n$ | $1 + 2\operatorname{sn}(\mathcal{E}_{\lambda}^{\otimes n}(\Psi))$ [Fig. 6] | $1 + \mathcal{O}\left(\left(\frac{d-1}{d}\right)^n\right)$ |

Take home messages

- Decay RoM of hypergraph states under independent noise.
- High-degree edges are fragile
- 3-complete hypergraph state
 - vanishing local magic (even on $\Theta(n)$ qubits)
 - non-vanishing magic threshold

Outlook

- Identify the most robust hypergraph states.
- The necessary and sufficient conditions for a family of states to have a non-vanishing magic threshold.
- The noise robustness of MBQC power.

Thank you for your attention!



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